Is it worth replacing 3DFgat by 4DVAR in CERA's ocean component?

Can we improve coupled consistency through data assimilation?

<u>Arthur Vidard</u> Florian Lemarié, Rémi Pellerej

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The original 4DVar problem:

$$J(\mathbf{x}) = \frac{1}{2} \left( \mathbf{x} - \mathbf{x}^{b} \right)^{T} \mathbf{B}^{-1} \left( \mathbf{x} - \mathbf{x}^{b} \right) + \frac{1}{2} \left( \mathcal{H}(\mathcal{M}(\mathbf{x})) - \mathbf{y}^{o} \right)^{T} \mathbf{R}^{-1} \left( \mathcal{H}(\mathcal{M}(\mathbf{x})) - \mathbf{y}^{o} \right)$$

It can be put in a more compact form, let  $F : \mathbb{R}^n \longrightarrow \mathbb{R}^{n+p}$  such that

$$F(\mathbf{x}) = \begin{pmatrix} \mathbf{B}^{-1/2} (\mathbf{x} - \mathbf{x}^b) \\ \mathbf{R}^{-1/2} (\mathcal{H}(\mathcal{M}(\mathbf{x})) - \mathbf{y}^o) \end{pmatrix}$$

Original cost function can then be rewritten

$$J(\mathbf{x}) = \frac{1}{2} \|F(\mathbf{x})\|_2^2$$

Denoting  $\mathbf{F}_{\mathbf{x}} = \begin{pmatrix} \mathbf{B}^{-1/2} \\ \mathbf{R}^{-1/2} \mathbf{H}_{\mathbf{x}} \mathbf{M}_{\mathbf{x}} \end{pmatrix}$  the jacobian (tangent linear) of F differentiated around  $\mathbf{x}$ , gradient and Hessian of J read

$$\nabla_{\mathbf{x}} J = \mathbf{F}_{\mathbf{x}}^{T} F(\mathbf{x}) \in \mathbb{R}^{n}$$
$$\nabla_{\mathbf{x}}^{2} J = \mathbf{F}_{\mathbf{x}}^{T} \mathbf{F}_{\mathbf{x}} + Q(\mathbf{x}) \in \mathbb{R}^{n \times n}$$

where  $Q(\mathbf{x})$  denotes the second order terms

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At each iterations

- Newton solves:  $\nabla^2_{\mathbf{x}^{(k)}}J\delta\mathbf{x}^{(k+1)}=-\nabla_{\mathbf{x}^{(k)}}J$
- Gauss-Newton solves:  $\mathbf{F}_{\mathbf{x}^{(k)}}^{T} \mathbf{F}_{\mathbf{x}^{(k)}} \delta \mathbf{x}^{(k+1)} = -\nabla_{\mathbf{x}^{(k)}} J$

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At each iterations

- Newton solves:  $\left(\mathbf{F}_{\mathbf{x}^{(k)}}^{\mathsf{T}}\mathbf{F}_{\mathbf{x}^{(k)}} + Q(\mathbf{x}^{(k)})\right)\delta\mathbf{x}^{(k+1)} = -\mathbf{F}_{\mathbf{x}^{(k)}}^{\mathsf{T}}F(\mathbf{x}^{(k)})$
- Gauss-Newton solves:  $\mathbf{F}_{\mathbf{x}^{(k)}}^{\mathsf{T}}\mathbf{F}_{\mathbf{x}^{(k)}}\delta\mathbf{x}^{(k+1)} = -\mathbf{F}_{\mathbf{x}^{(k)}}^{\mathsf{T}}\mathcal{F}(\mathbf{x}^{(k)})$

At each iterations

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Under several conditions, Gauss Newton will converge toward a minimum of the original problem if  $\exists \eta^{(k)} < 1$ :

$$\left\| \left( \mathsf{F}_{\mathsf{x}^{(k)}}^{\mathsf{T}} \mathsf{F}_{\mathsf{x}^{(k)}} \right)^{-1} Q(\mathsf{x}^{(k)}) \right\|_{2} \leq \eta^{(k)}$$

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In practice further approximations are made (lower resolution, simplified physics, CERA, 3DFgat, ...), the approximate Gauss-Newton iteration then solves

$$\widetilde{\mathbf{F}}_{\mathbf{x}^{(k)}}^{\mathsf{T}}\widetilde{\mathbf{F}}_{\mathbf{x}^{(k)}}\delta\mathbf{x}^{(k+1)} = -\widetilde{\mathbf{F}}_{\mathbf{x}^{(k)}}^{\mathsf{T}}F(\mathbf{x}^{(k)})$$

One can show that for such an approximation of the cost function, this sufficient condition becomes

$$\left\| \mathbf{I} - \left( \widetilde{\mathbf{F}}_{\mathbf{x}^{(k)}}^{\mathcal{T}} \widetilde{\mathbf{F}}_{\mathbf{x}^{(k)}} \right)^{-1} \left( \widetilde{\mathbf{F}}_{\mathbf{x}^{(k)}}^{\mathcal{T}} \mathbf{F}_{\mathbf{x}^{(k)}} + \widetilde{Q}(\mathbf{x}^{(k)}) \right) \right\|_{2} \leq \eta^{(k)}$$

But the minimum is not the same as the original problem

$$\|\tilde{\mathbf{x}}^* - \mathbf{x}^*\|_2 \leq \frac{1}{1-\nu} \left\| \left( \widetilde{\mathbf{F}}_{\tilde{\mathbf{x}}^*}^+ - \mathbf{F}_{\tilde{\mathbf{x}}^*}^+ \right) F(\tilde{\mathbf{x}}^*) \right\|_2 = \frac{1}{1-\nu} \left\| \mathbf{F}_{\tilde{\mathbf{x}}^*}^+ F(\tilde{\mathbf{x}}^*) \right\|_2$$

 $(\mathbf{F}^{+} = \left(\mathbf{F}^{\mathcal{T}}\mathbf{F}\right)^{-1}\mathbf{F}^{\mathcal{T}})$ 

In the linear case the above sufficient condition becomes necessary.

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 $(\mathbf{F}^{+} = \left(\mathbf{F}^{\mathcal{T}}\mathbf{F}\right)^{-1}\mathbf{F}^{\mathcal{T}})$ 

In the linear case the above sufficient condition becomes necessary.

As a summary all what matters is:

how good 
$$\widetilde{F}_x = \begin{pmatrix} B^{-1/2} \\ R^{-1/2} H_x \widetilde{M}_x \end{pmatrix}$$
 is an approximation of  $F_x = \begin{pmatrix} B^{-1/2} \\ R^{-1/2} H_x M_x \end{pmatrix}$ .

#### Back to the first question Is it worth replacing 3DFgat by 4DVAR in CERA20C's ocean component?

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#### Back to the first question Is it worth replacing 3DFgat by 4DVAR in CERA20C's ocean component?

the answer is no ...

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#### Back to the first question Is it worth replacing 3DFgat by 4DVAR in CERA20C's ocean component?

the answer is no ...

It just does not change a bit ...

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ORCA 1, One day assimilation window, T and S assimilation



4Dvar - 3Dvar (1 day) difference T300 increment



ORCA 1, 30 day assimilation window, T and S assimilation







Considering the standard inner loop's incremental formulation:

$$J^{k}(\delta \mathbf{x}^{(k)}) = \left( \delta \mathbf{x}^{(k)} + \sum_{l=1}^{(k-1)} \delta \mathbf{x}^{(l)} \right)^{T} \mathbf{B}^{-1} \left( \delta \mathbf{x}^{(k)} + \sum_{l=1}^{(k-1)} \delta \mathbf{x}^{(l)} \right) \\ + \sum_{i=0}^{N} \left( \mathbf{H}_{t_{i}}^{(k-1)} \tilde{\mathbf{M}}_{t_{i}}^{(k-1)} \delta \mathbf{x}^{(k)} - \mathbf{d}_{t_{i}}^{(k-1)} \right)^{T} \mathbf{R}_{t_{i}}^{-1} \left( \mathbf{H}_{t_{i}}^{(k-1)} \tilde{\mathbf{M}}_{t_{i}}^{(k-1)} \delta \mathbf{x}^{(k)} - \mathbf{d}_{t_{i}}^{(k-1)} \right)$$

How good our approximation of the "true" F (i.e. HM) is? For 3D-Var ( $\tilde{M} = I$ ) and 4D-Var ( $\tilde{M}$  includes some approximations as well)?

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Approximation in the linear propagator

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Approximation in the linear propagator

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ORCA 025, One day assimilation window, T,S and SSH assimilation



4Dvar - 3Dvar (1 day) T300 increment difference



ORCA 025, 5 day assimilation window, T,S and SSH assimilation



4Dvar - 3Dvar (5 days) T300 increment difference



There are potential interest to use 4DVar for longer assimilation windows / higher resolution, but it comes at a cost:

Orca1, 10 iteration, 1 node:

1day:

4dvar: 12mn (17mn) 3dvar: 6mn (11mn)

Orca025, 5 iteration, 6 nodes:

5 day:

4dvar: 7h (9h) 3dvar: 45mn (2h45) 10days:

4dvar: 48mn (1h) 3dvar: 6mn (16mn)

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### Simplified 4Dvar Do we really need a full tangent model?

$$\frac{\partial \delta T}{\partial t} = -\nabla . (\delta T \mathbf{U}) + \delta D^{vT}$$
$$\delta D^{vT} = \frac{\partial}{\partial z} \left( A^{vT} \frac{\partial \delta T}{\partial z} \right)$$



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# Simplified 4Dvar

Single temperature observation (10d assimilation window)



### Multi incremental 4Dvar

Do we really need a full resolution?

ORCA025 for the direct model, ORCA1 for the tangent model. Perturbations generated at coarse resolution.

Interpolation: observation operator

simplification: its weighted ajoint.



#### Approximation in the linear propagator

## Multi incremental 4Dvar



4dvar: 45mn 3dvar: 2mn

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3dvar: 45mn (2h45)

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Back to the second question Can we improve coupled consistency through data assimilation?

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#### Back to the second question Can we improve coupled consistency through data assimilation?

the answer is yes, probably ...

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#### Back to the second question Can we improve coupled consistency through data assimilation?

the answer is yes, probably ...

but at a cost ...

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OA coupling is a complex matter with many sources of uncertainties

- time/space non-confomity
- interfaces may actually not be represented by any component
- multi physics with different characteristics.
- highly parameterised interface (Bulk formulae)
- coupling methods

Some of theses uncertainties are unavoidable, some others are linked to the way we implement things.

Coupled DA is an opportunity to account for or reduce them

<sup>• ...</sup> 

## Focus on flux consistency



The SWR algorithm reads :

$$\begin{cases} \mathcal{L}_{a}(u_{a}^{k}) = f_{a} & \text{on } \Omega_{a} \times T_{W} \\ u_{a}^{k}(z,0) = u_{0}(z) & z \in \Omega_{a} \\ \mathcal{C}_{a}(u_{a}^{k}) = \mathcal{F}_{oa}(u_{a}^{k}, u_{o}^{k-1}) & \text{on } \Gamma \times T_{W} \end{cases} \begin{cases} \mathcal{L}_{o}(u_{o}^{k}) = f_{o} & \text{on } \Omega_{o} \times T_{W} \\ u_{o}^{k}(z,0) = u_{0}(z) & z \in \Omega_{o} \\ \mathcal{C}_{o}(u_{o}^{k}) = \mathcal{F}_{oa}(u_{a}^{k}, u_{o}^{k}) & \text{on } \Gamma \times T_{W} \end{cases}$$
  
where  $T_{W} = [t_{i}; t_{i+1}]$ 

• At convergence, it provides a flux consistent solution :  $C_a(u_a) = C_o(u_o)$  on  $\Gamma \times T_W$ 

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### For a fistful of algorithms

• Fully coupled models.  $\mathbf{x} = u_0(z), z \in \Omega$ 

$$J_{FCM}(\mathbf{x}) = \left(\mathbf{x} - \mathbf{x}^{b}\right)^{T} \mathbf{B}^{-1} \left(\mathbf{x} - \mathbf{x}^{b}\right) + \sum_{i=0}^{N} \left(\mathcal{H}_{t_{i}}(\mathcal{M}_{t_{i}}(\mathbf{x})) - \mathbf{y}_{t_{i}}^{o}\right)^{T} \mathbf{R}_{t_{i}}^{-1} \left(\mathcal{H}_{t_{i}}(\mathcal{M}_{t_{i}}(\mathbf{x})) - \mathbf{y}_{t_{i}}^{o}\right)$$

• Partially coupled models.  $\mathbf{x}_0 = (u_0(z), u_o^0(0, t))^T$ ,  $z \in \Omega$ ,  $t \in [0, T]$ 

$$J_{PCM}(\mathbf{x}) = J^{b}(\mathbf{x}) + \sum_{i=0}^{N} \left( \mathcal{H}_{t_{i}}(\mathcal{M}_{t_{i}}^{trunc}(\mathbf{x})) - \mathbf{y}_{t_{i}}^{o} \right)^{T} \mathbf{R}_{t_{i}}^{-1} \left( \mathcal{H}_{t_{i}}(\mathcal{M}_{t_{i}}^{trunc}(\mathbf{x})) - \mathbf{y}_{t_{i}}^{o} \right) + J^{s}(\mathbf{x})$$

- Weakly coupled models.  $\mathbf{x}_0 = (u_0(z), u_o^0(0, t), u_a^0(0, t))^T, z \in \Omega, t \in [0, T]$  $J_{WCM}(\mathbf{x}) = J_a^b(\mathbf{x}_a) + J_o^b(\mathbf{x}_o) + J_a^o(\mathbf{x}_a) + J_o^o(\mathbf{x}_o) + J^s(\mathbf{x})$
- and obviously CERA (uncoupled in the inner loop).  $\mathbf{x} = u_0(z), \ z \in \Omega$

$$J^{\mathfrak{s}}(\mathbf{x}) = \gamma \|\mathcal{C}_{\mathfrak{a}}(u_{\mathfrak{a}}(0,t)) - \mathcal{C}_{\mathfrak{o}}(u_{\mathfrak{o}}(0,t))\|_{[0,T]}^2$$

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#### Stand-alone SCM



Figure: Colors represent zonal (a,c) and meridional (b,d) atmosphere wind and ocean current velocities components and black isolines represent temperatures.

$$\frac{\partial \mathbf{u}_{\beta}(z,t)}{\partial t} = -f\mathbf{k} \times \mathbf{u}_{\beta}(z,t) + \frac{\partial}{\partial z} \left( K_{m}^{\beta}(z) \frac{\partial \mathbf{u}_{\beta}(z,t)}{\partial z} \right) + F_{\mathbf{u}_{\beta}}(z,t) \quad \text{sur } \Omega_{\beta} \times [0,T]$$
$$\frac{\partial \mathbf{t}_{\beta}(z,t)}{\partial t} = \frac{\partial}{\partial z} \left( K_{s}^{\beta}(z) \frac{\partial \mathbf{t}_{\beta}(z,t)}{\partial z} \right) + F_{\mathbf{t}_{\beta}}(z,t) \quad \text{sur } \Omega_{\beta} \times [0,T]$$

where  $\beta = a, o$  refer to atmosphere and ocean variables respectively. Both models use the same structure and differ from their forcing terms  $F_*$ , their interface conditions and the computation of their turbulent viscosity and diffusivity coefficients  $K_m^\beta$  and  $K_s^\beta$ .

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Algorithm	$\gamma$	k <sub>max</sub>	# of minimisation	Computing cost	Interface imbalance	RMSE improvement
			iterations	(relative to CERA)	indicator	(in %)
FCM-F	-	kcvg	26	3.8	2.10-12	74
CERA-F	-	k <sub>cvg</sub>	24	1.1	5.810-12	24
CERA-1	-	1	26	1	1.6	40
CERA-1-SWR	-	1	26	1	5 10 <sup>-2</sup>	45
PCM-1	0.1	1	25	0.96	4.10 <sup>-3</sup>	60
WCM	0.1	0	31	1.2	6.10-3	57

Table: Result summary for the SCM system (limited to 2 outer loops)



Figure: Forecast of SSU and SSV from FCM-F, CERA-1 and WCM analysis. Dashed and plain black lines are background and truth evolutions respectively

In the previous frame, we were limited to 2 outer loops, due to non convergence of CERA. Adding the  $J^s$  term sorts this out



In the previous frame, we were limited to 2 outer loops, due to non convergence of CERA. Adding the  $J^s$  term sorts this out But we can inflate **B** to make CERA converge



The outer/inner loop framework allows for approximation

- their impact can be studied theoretically
- they can (should?) be specific for a given application
- they can be (partially) accounted for by modifying the cost function

In addition to deliverables

- 4DVar, simplistic 4DVar and multigrid 4DVar are available in Nemovar repository
- a stand alone single column will soon be available along with its OOPs interfaces

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